Lattice QCD for Nuclear Physics

- 1. Lattice QCD infrastructure: people and software
- 2. Hadronic Parity Violation (PV): I=2 NN PV Amplitude
- 3. Neutrinoless Double Beta-Decay (0νββ)
- 4. Nucleon Matrix Elements for Fundamental Symmetry Tests

Tight-knit team

red = postdoc blue = grad student

Lattice QCD Team

LBL/UCB: David Brantley, Chia Cheng (Jason) Chang, W. Haxton, T. Kurth (NERSC), Ken McElvain, Henry Monge Camacho, Amy Nicholson, AWL

LLNL: Evan Berkowitz, Enrico Rinaldi, Pavlos Vranas

Liverpool/Plymouth: Nicolas Garron

JLab: Balint Jóo

CCNY/RIKEN: Brian Tiburzi

NVIDIA: Kate Clark

Computational Scientists

LBL: Fernando Perez, Abhinav Sarje, Sam Williams

LBL (seeking \$): Pieter Ghysels, Khaled Ibrahim









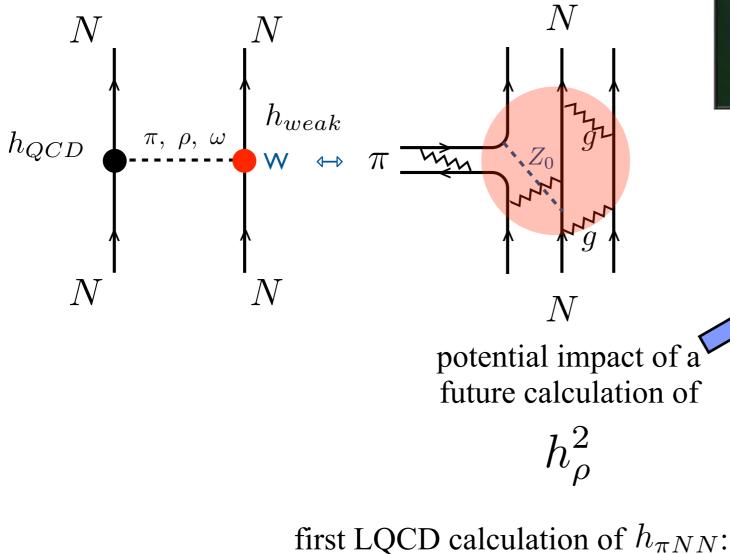


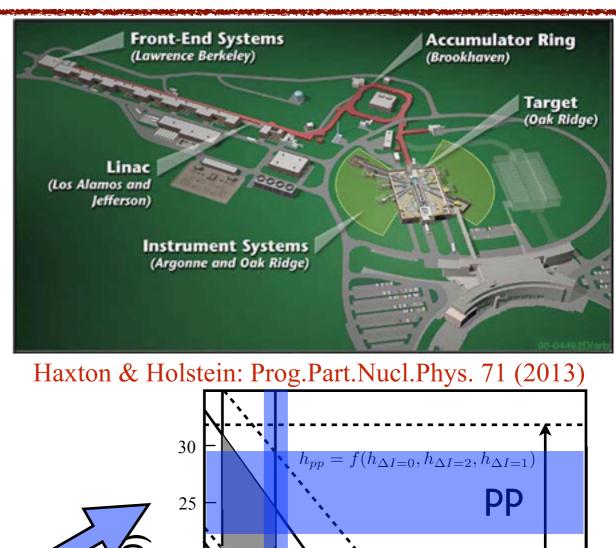
Lattice QCD (LQCD) Software

- Our main software, latscat, was initially developed by Thorsten Kurth (now @ NERSC). Evan Berkowitz, Enrico Rinaldi and AWL are now co-developers with Thorsten consulting regularly
- We work closely with Kate Clark (NVIDIA) who develops QUDA (GPU library) and continues to optimize software for our projects (Titan/OLCF)
- We work closely with Balint Joo (principal developer of USQCD libraries)
 who continues to optimize and add support for routines we use
- We work closely with Abhinav Sarje (LBL CRD through CalLat SciDAC3) and Ken McElvain (20+ yr software engineer turned UC Berkeley physics grad student) on significant performance optimizations, O(50-100%)
- We have implemented a **NERSC database** to share numerical results amongst ourselves and ultimately, to share our raw LQCD results and our analysis results (the physics) publicly with easy open access to interested physicists
- We are seeking funding for partial support of LBL Computer Scientist/Applied Math scientists who are familiar with exascale development and LQCD to develop next-generation LQCD code for NP on (near-)exascale computers

Hadronic Parity Violation (HPV

• The first CalLat lattice QCD project planned was to compute the Isospin=2 Hadronic Parity Violating Amplitude between 2 Nucleons





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 $L=2.5 \text{ f a}=0.123 \text{ f m}_{\pi}=389 \text{ MeV}$ systematic approximations J. Wasem Phys. Rev. C85 (2012) 022501 10 |18**|** 1-6-L pO 10 12 $h(\Delta I = I)$

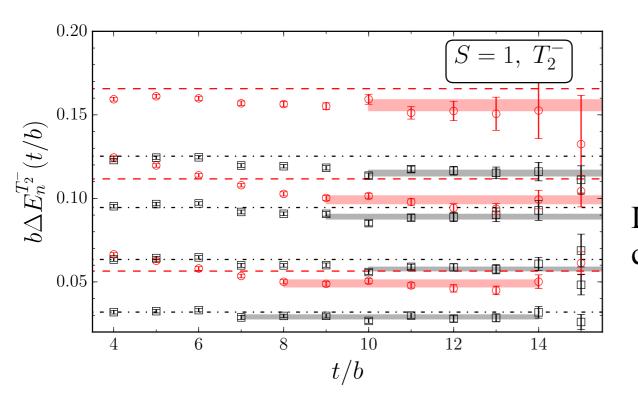
Hadronic Parity Violation (HPV)

• We needed to develop new technology to compute higher partial wave NN scattering phase shifts (S,P,D,F), displaced two-nucleon interpolating fields $N^{\dagger}(t_0, x_0) N^{\dagger}(t_0, x_0 + r_0) | 0 >$

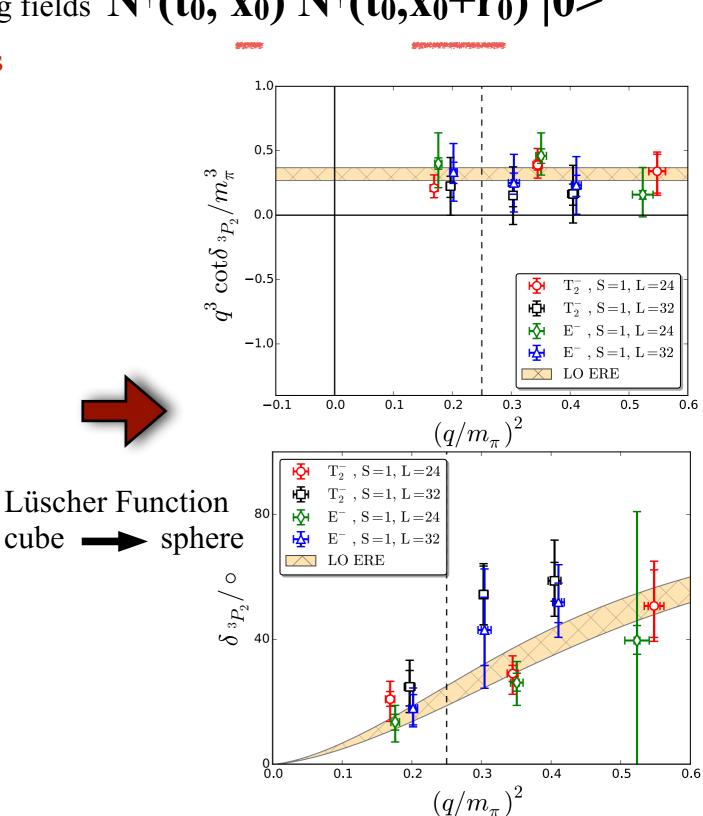
First LQCD calculation of NN P, D (,F) waves arXiv:1508.00886 $(m_\pi \simeq 800 \ {\rm MeV})$

Example: ³P₂ - elastic wave

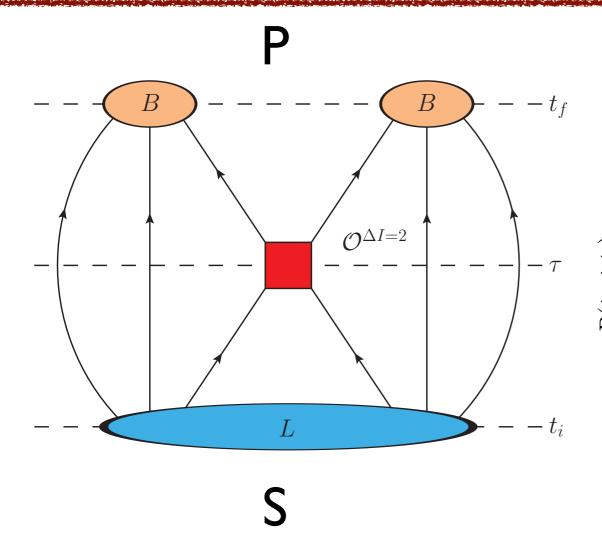
Lattice QCD calculation of NN energy levels



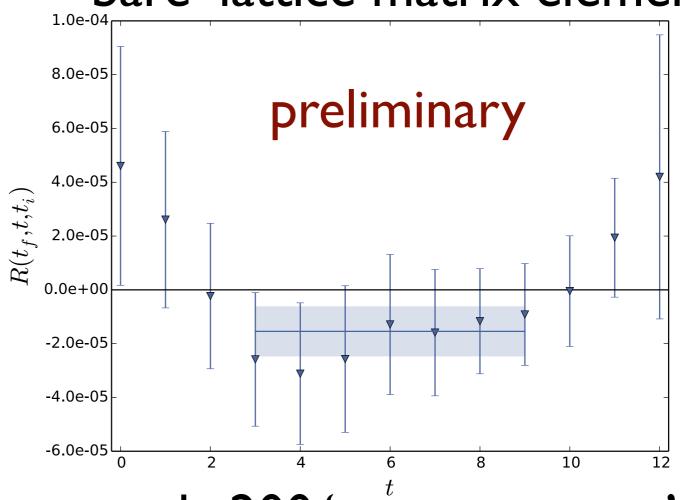




Hadronic Parity Violation (HPV)



'bare' lattice matrix element



only 200 'measurements'

$$\langle pp(^{3}P_{1})|\mathcal{O}^{\Delta I=2}|pp(^{1}S_{0})\rangle_{V=\infty} = LL\left(\delta_{^{1}S_{0}}, \frac{\partial\delta_{^{1}S_{0}}}{\partial E}, \delta_{^{3}P_{1}}, \frac{\partial\delta_{^{3}P_{1}}}{\partial E}\right)\langle pp(^{3}P_{1})|\mathcal{O}^{\Delta I=2}|pp(^{1}S_{0})\rangle_{V=L^{3}}$$

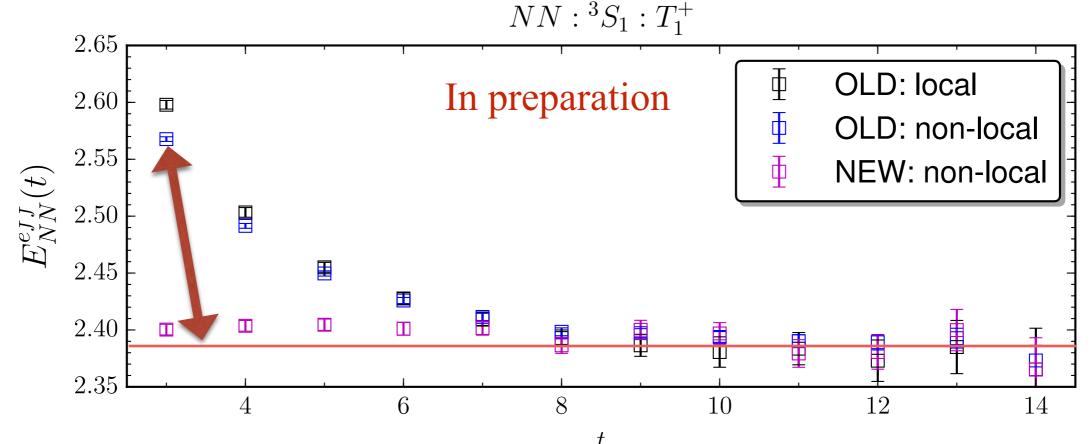
Known Lellouch-Lüscher (LL) function

Raúl Briceño et al

Most significant systematic/challenge for HPV is our determination of good 2-nucleon operators.

Further Improved NN Operators

• We have developed new operators which may allow for "exponentially" improved results.



= our old determination of the ground state (gs) energy with OLD operators



= deviation from gs plateau is driven by excited state contamination

NEW: non-local

= our new strategy - the excited state contamination of the two-nucleon correlation function has been significantly reduced. This will allow analysis beginning earlier in time where the stochastic noise is ~ exponentially smaller

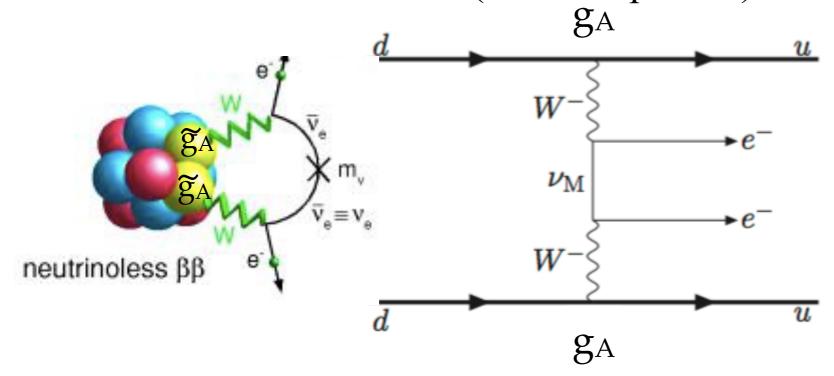
OLD: non-local

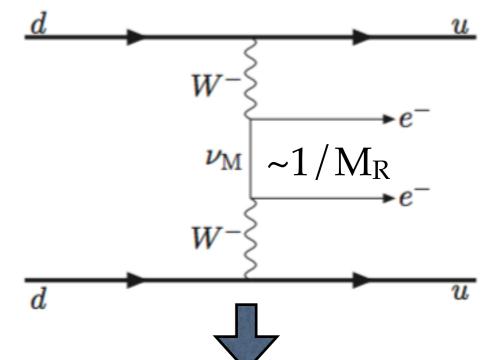
= our recently developed displaced two-nucleon interpolating fields (see previous slides) $N^{\dagger}(t_0, x_0) N^{\dagger}(t_0, x_0 + r_0) | 0 >$

OLD: local = two-nucleon interpolating operators from the same space-time location - the strategy used by most other groups, $N^{\dagger}(t_0, x_0) N^{\dagger}(t_0, x_0) | 0 >$

long-range contribution (standard picture)

short-range contribution possibly equally/more important





Long Range: lattice QCD can help understand "quenching" of g_A in a nucleus

Short Range: lattice QCD is the ONLY theoretical tool we have to understand these contributions with quantified uncertainties

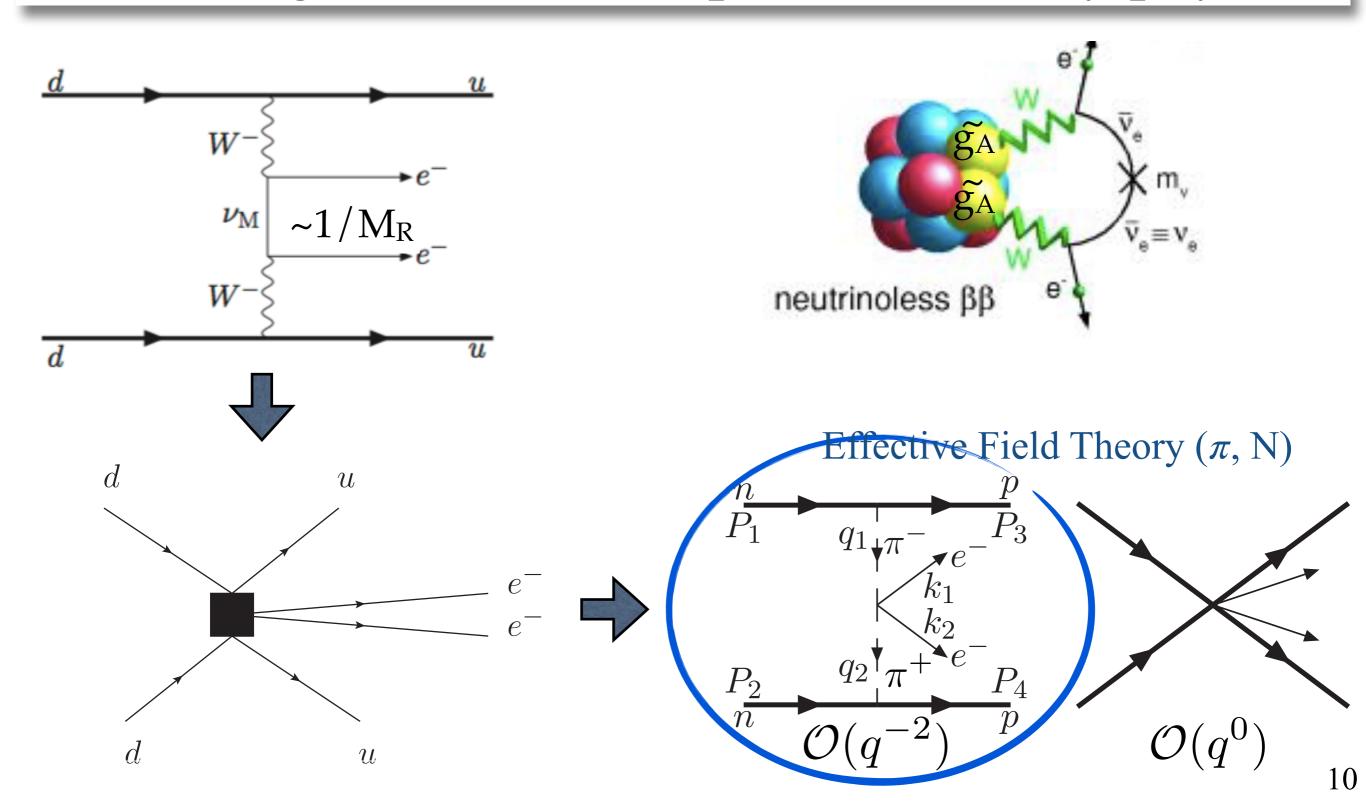
Lattice QCD: compute 2-nucleon matrix elements to determine unknown couplings/transition rates

Many Body Nuclear Effective Theory: take lattice QCD results as input and compute transition rate in nucleus (Haxton, others)

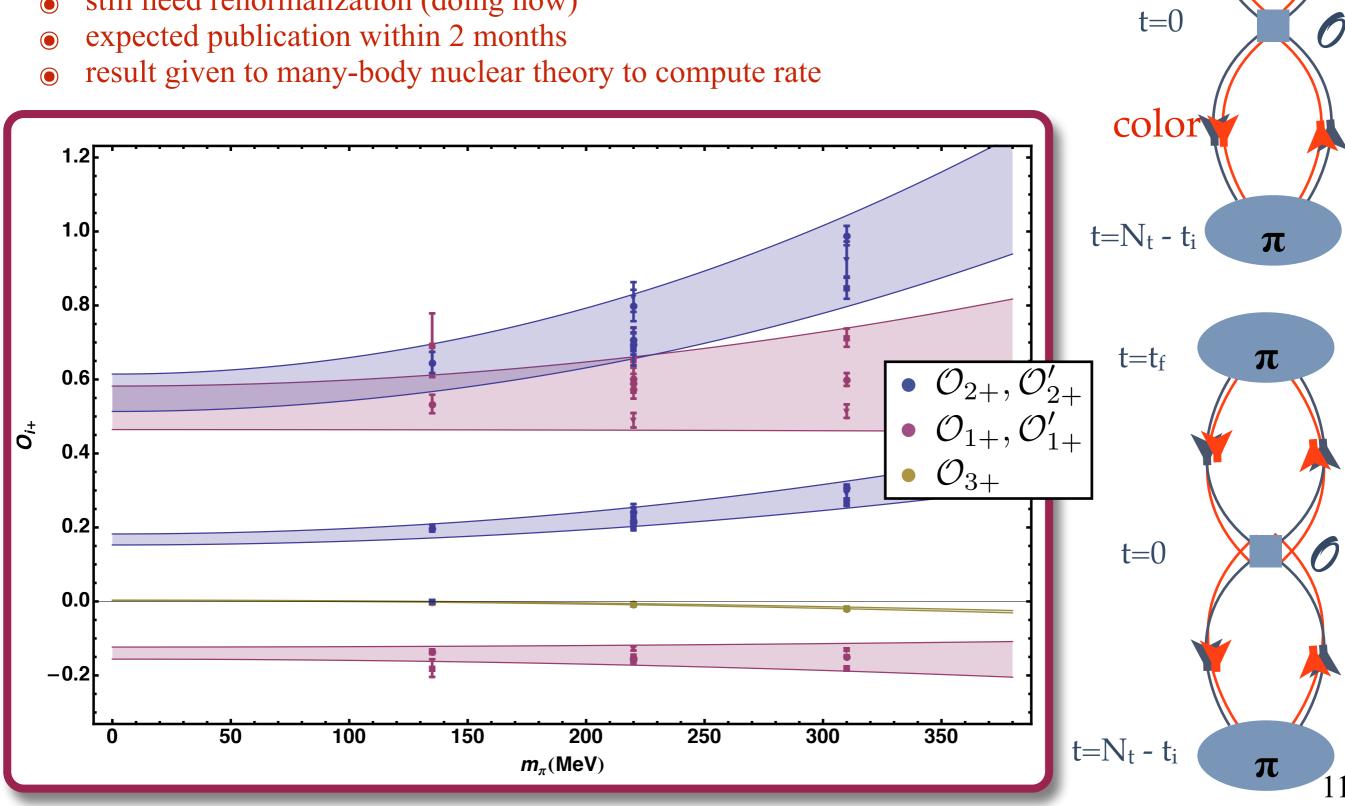
In this review, focus on short-range contributions

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Short-range contribution: probe for heavy physics



We have performed the first, AND COMPLETE lattice QCD calculation of the $\pi^- \to \pi^+$ transition amplitude which is expected to dominate the $0\nu\beta\beta$ rate in the case of short distance physics $\pi^- \to \pi^+$ physical pion mass ocontinuum and infinite volume limits still need renormalization (doing now) expected publication within 2 months



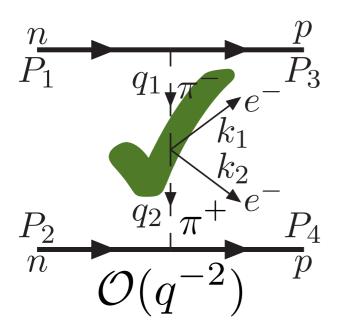
 $t=t_f$

spin

π

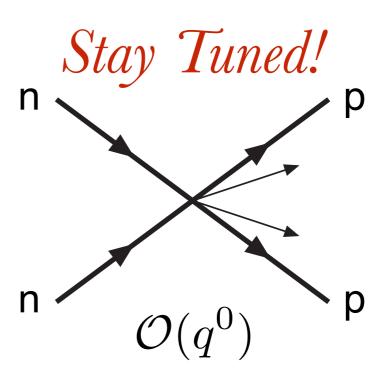
Short-distance Contact operators

- LO almost complete!
 - dominant contribution to 0vββ from short-range contributions
 - need renormalization



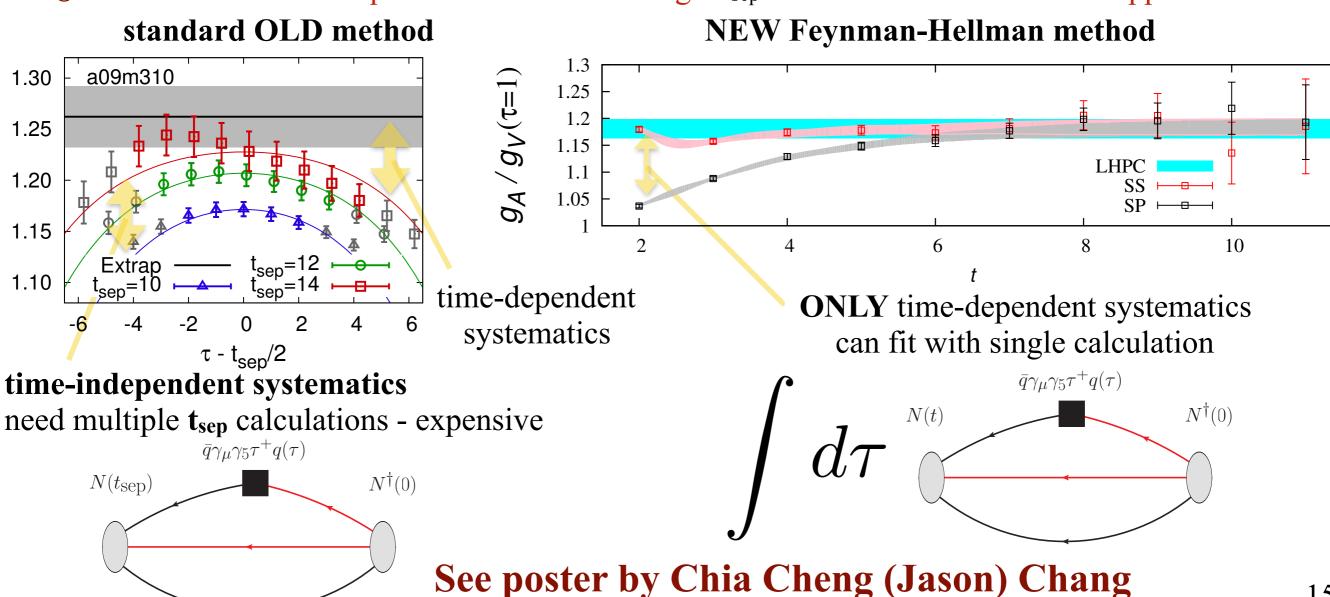
$$V_{nn\to pp}(q_1,q_2)$$

•nn → pp contact operators
 next step: similar to two-N
 hadronic parity violation



- Low energy precision tests of the Standard Model place competitive bounds on physics beyond the Standard Model (BSM) as compared to the LHC.
- Interpreting null results/hopeful signals in terms of potential BSM physics requires a quantitative and often precise knowledge of one- and sometimes two-nucleon matrix elements
 - direct dark matter detection
 - permanent electric dipole moments in nucleons and nuclei
 - \bullet modification from V-A weak beta decay, $n \to p + e^- + \bar{\nu}_e$
 - \bullet $\mu \rightarrow$ e conversion
 - ...
- The largest uncertainty in many examples comes from the hadronic uncertainty of the nucleon matrix elements
- Lattice QCD calculations of nucleon matrix elements have one additional systematic that is more complicated than regular spectrum calculations: excited state systematics

- With C. Bouchard, K. Orginos, we have developed a new method for computing matrix elements
 - significantly improved control over excited states systematics, based upon the Feynman-Hellman Theorem, in which all excited states contributions are time-dependent (controlled in fit to numerical results)
 - ~9 times more statistical results for approximately equal computing time
 - \odot entire calculation is equivalent in cost to 1 single t_{sep} calculation of the standard approach



Embargoed Result

Lattice QCD for Nuclear Physics

- 1. Lattice QCD infrastructure: people and software
 - a) we have a very active and tight-knit/engaged team of physicists and computer scientists
 - b) we are developing independent software for our specialized needs
- 2. Hadronic Parity Violation (PV): I=2 NN PV Amplitude
 - a) PV calculation on pause as we need to improve two-nucleon elastic scattering calculations
 - b) we have developed an exciting new method that may provide exponential improvement in NN LQCD calculations publication coming soon
- 3. Neutrinoless Double Beta-Decay $(0\nu\beta\beta)$
 - a) 0νββ may receive important contributions from short-range 4-quark—2-electron operators (more general BSM theory causing Lepton-number violation)
 - b) The term expected to dominate such a contribution comes from a short-range Isospin=2 pion matrix element $\pi^- \to \pi^+$
 - c) We have nearly completed the calculation of this matrix element with LQCD all that remains is the non-perturbative renormalization which we are performing now
 - d) next is the calculation of two-nucleon operators same technology as HPV
- 4. Nucleon Matrix Elements for Fundamental Symmetry Tests
 - a) developed new method for computing matrix elements motivated by the Feynman-Hellman Theorem
 - b) applied this new method to calculate the nucleon axial charge